

Normal and Anomalous Fluctuation Relations for Gaussian Stochastic Dynamics

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Motivation: Fluctuation relations

Consider a (classical) particle system evolving from some initial state into a nonequilibrium steady state.

Measure the probability distribution $\rho(\xi_t)$ of entropy production ξ_t during time t :

$$\ln \frac{\rho(\xi_t)}{\rho(-\xi_t)} = \xi_t$$

Transient Fluctuation Relation (TFR)

Evans, Cohen, Morriss (1993); Gallavotti, Cohen (1995)

why important? of *very general validity* and

- 1 generalizes the **Second Law** to (small) systems in nonequ.
- 2 connection with **fluctuation dissipation relations**
- 3 can be checked in **experiments** (Wang et al., 2002)

Fluctuation relation for Langevin dynamics

warmup: check TFR for the **overdamped Langevin equation**

$$\dot{x} = F + \zeta(t) \quad (\text{set all irrelevant constants to 1})$$

for a particle at position x with constant field F and noise ζ .

entropy production ξ_t is equal to (mechanical) **work** $W_t = Fx(t)$

with $\rho(W_t) = F^{-1} \varrho(x, t)$; remains to solve corresponding

Fokker-Planck equation for initial condition $x(0) = 0$:

the position pdf is Gaussian,

$$\varrho(x, t) = \frac{1}{\sqrt{2\pi\sigma_x^2}} \exp\left(-\frac{(x - \langle x \rangle)^2}{2\sigma_x^2}\right)$$

straightforward:

$$\text{(work) TFR holds if } \langle x \rangle = \sigma_x^2/2$$

and \exists **fluctuation-dissipation relation 1 (FDR1) \Rightarrow TFR**

see, e.g., **van Zon, Cohen, PRE (2003)**

Gaussian stochastic dynamics

goal: check TFR for **Gaussian stochastic processes** defined by the (overdamped) **generalized Langevin equation**

$$\int_0^t dt' \dot{x}(t') K(t-t') = F + \zeta(t)$$

e.g., **Kubo (1965)**

with **Gaussian noise** $\zeta(t)$ and **memory kernel** $K(t)$

such dynamics can generate **anomalous diffusion**:

$$\sigma_x^2 \sim t^\alpha \text{ with } \alpha \neq 1 (t \rightarrow \infty)$$

TFR for correlated internal Gaussian noise

consider two generic cases:

1. **internal Gaussian noise** defined by the **FDR2**,

$$\langle \zeta(t)\zeta(t') \rangle \sim K(t-t'),$$

with **non-Markovian (correlated) noise**; e.g., $K(t) \sim t^{-\beta}$

solving the corresponding generalized Langevin equation in Laplace space yields

$$\text{FDR2} \Rightarrow \text{'FDR1'}$$

and since $\rho(W_t) \sim \varrho(x, t)$ is Gaussian

$$\text{'FDR1'} \Rightarrow \text{TFR}$$

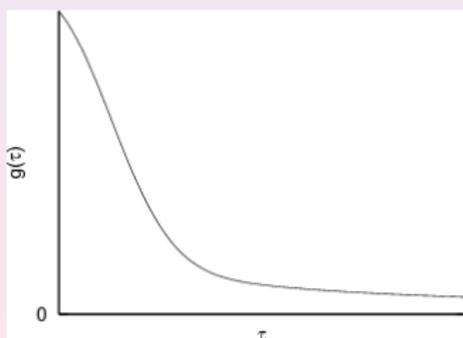
for correlated internal Gaussian noise \exists TFR

Correlated external Gaussian noise

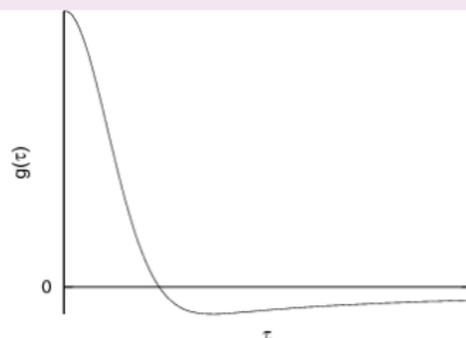
2. **external Gaussian noise** for which there is **no FDR2**, modeled by the (overdamped) generalized Langevin equation

$$\dot{x} = F + \zeta(t)$$

consider two types of **Gaussian noise correlated** by $g(\tau) = \langle \zeta(t)\zeta(t') \rangle_{\tau=t-t'} \sim (\Delta/\tau)^\beta$ for $\tau > \Delta$, $\beta > 0$:



persistent



antipersistent

it is $\langle x \rangle = Ft$ and $\sigma_x^2 = 2 \int_0^t d\tau (t - \tau)g(\tau)$

Results: Anomalous TFRs

σ_X^2 and the **fluctuation ratio** $R(W_t) = \ln \frac{\rho(W_t)}{\rho(-W_t)}$ for $t \gg \Delta$ and $g(\tau) = \langle \zeta(t)\zeta(t') \rangle_{\tau=t-t'} \sim (\Delta/\tau)^\beta$:

β	persistent		antipersistent *	
	σ_X^2	$R(W_t)$	σ_X^2	$R(W_t)$
$0 < \beta < 1$	$\sim t^{2-\beta}$	$\sim \frac{W_t}{t^{1-\beta}}$	regime does not exist	
$\beta = 1$	$\sim t \ln\left(\frac{t}{\Delta}\right)$	$\sim \frac{W_t}{\ln\left(\frac{t}{\Delta}\right)}$		
$1 < \beta < 2$			$\sim t^{2-\beta}$	$\sim t^{\beta-1} W_t$
$\beta = 2$	$\sim 2Dt$	$\sim \frac{W_t}{D}$	$\sim \ln(t/\Delta)$	$\sim \frac{t}{\ln\left(\frac{t}{\Delta}\right)} W_t$
$2 < \beta < \infty$			$= \text{const.}$	$\sim t W_t$

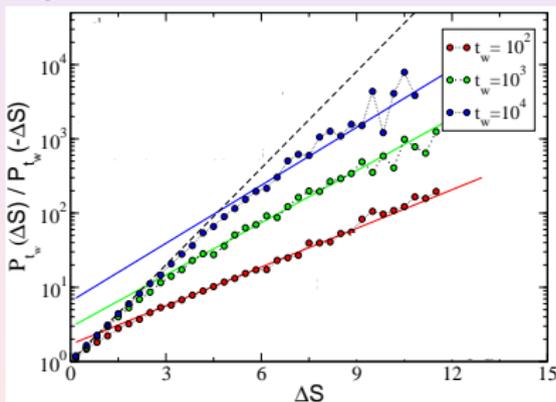
* antipersistence for $\int_0^\infty d\tau g(\tau) > 0$ yields **normal diffusion** with **generalized TFR**; above antipersistence for $\int_0^\infty d\tau g(\tau) = 0$

Relations to experiments: glassy dynamics

$$R(W_t) = \ln \frac{\rho(W_t)}{\rho(-W_t)} = \mathbf{f}_\beta(\mathbf{t}) W_t$$

means by plotting R for different t the **slope might change**.

example 1: computer simulations for a **binary Lennard-Jones mixture** below the glass transition

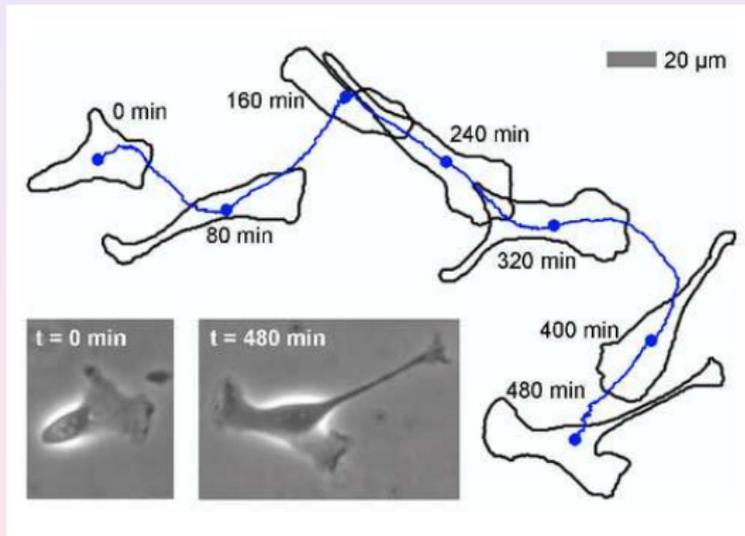


Crisanti, Ritort, PRL (2013)

- similar results for other glassy systems (Sellitto, PRE, 2009)

Biological cell migration

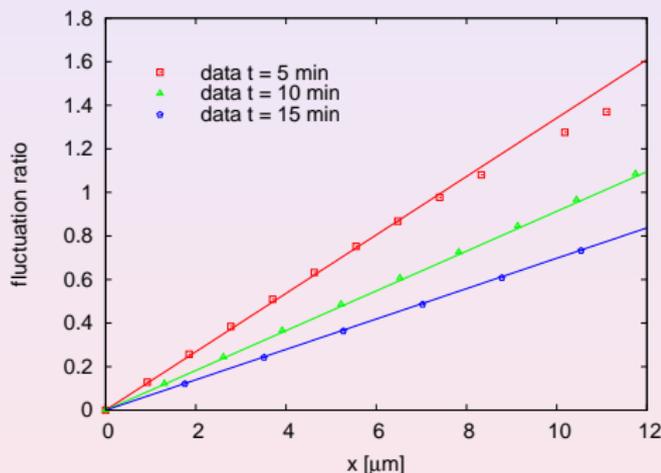
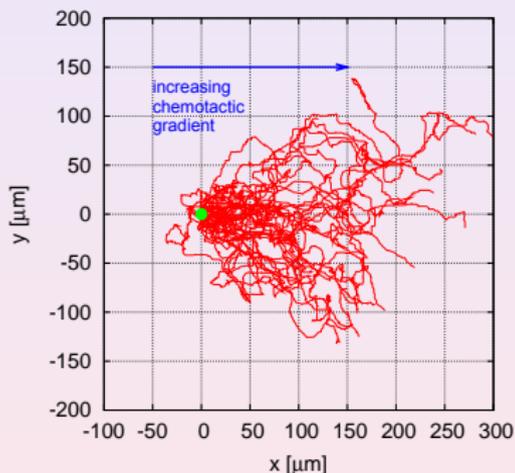
example 2: **single biological cell** crawling on a substrate; trajectory recorded with a video camera



Dieterich, RK et al., PNAS, 2008

Cell migration under chemical gradients

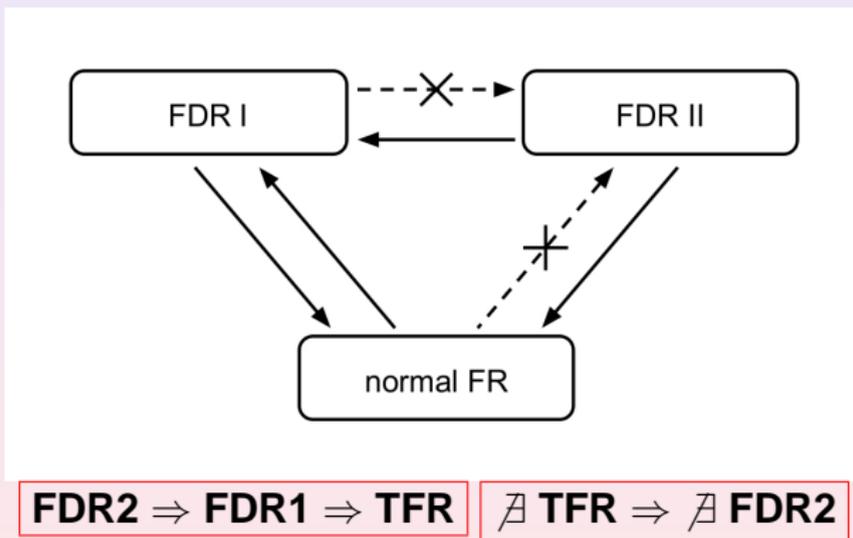
experiments on **murine neutrophils** under **chemotaxis**:



Dieterich et al. (2013)

Summary

- **relation between TFR and FDR I,II** for correlated Gaussian stochastic dynamics: ('normal FR'= conventional TFR)



- **anomalous TFRs** likely to be important for **glassy dynamics**

References

- A.V. Chechkin, F.Lenz, RK, J. Stat. Mech. L11001 (2012)
- RK, A.V. Chechkin, P.Dieterich, *Anomalous fluctuation relations* in:

