Diffusion in randomly perturbed dissipative dynamics

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Outline

Motivation

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- Motivation: standard map, diffusion, and dissipation
- Randomly perturbed dissipative dynamics
- Randomly perturbed dissipative standard map
- Continuous Time Random Walk theory

• paradigmatic Hamiltonian dynamical system:

standard map

$$x_{n+1} = x_n + y_n \mod 2\pi$$

$$y_{n+1} = y_n + K \sin x_{n+1}$$

derived from kicked rot(at)or where $x_n \in \mathbb{R}$ is an angle, $y_n \in \mathbb{R}$ the angular velocity with $n \in \mathbb{N}$ and K > 0 the kick strength

define diffusion coefficient as

$$D(K) = \lim_{n \to \infty} \frac{1}{n} < (y_n - y_0)^2 >$$

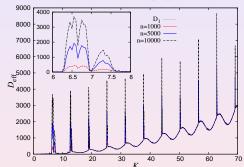
with ensemble average over the initial density

$$<...>=\int dx dy \ \varrho(x,y)..., x \in [0,2\pi), \ y=y_0 \in [0,2\pi)$$

Diffusion in the standard map

Outline

analytical (Rechester, White, 1980) and numerical studies of parameter-dependent diffusion $D_{eff}(K)$:



Manos, Robnik, PRE (2014)

- D(K) is highly irregular
- for some K there is superdiffusion with mean square displacement $< y_n^2 > \sim n^{\gamma}$, $\gamma > 1$ due to accelerator modes

Summary

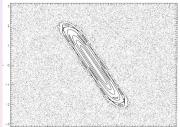
The dissipative standard map

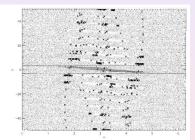
model damping in the standard map by

$$x_{n+1} = x_n + y_n \mod 2\pi$$

 $y_{n+1} = (1 - \nu)y_n + f_0 \sin x_{n+1}$

with $\nu \in [0, 1]$:





Feudel, Grebogi, Hunt, Yorke, PRE (1996)

- islands in phase space for $\nu = 0$ (left) become coexisting periodic attractors (right): 150 found for $\nu = 0.02$, $f_0 = 4$
- simple argument yields $|y_n| < y_{max}$: quick trapping

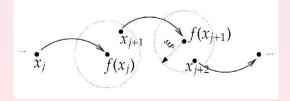
Question: What happens to dissipative deterministic dynamics $\mathbf{x}_{n+1} = \mathbf{f}(\mathbf{x}_n)$ under random perturbations?

Consider the dissipative standard map with additive noise:

$$x_{n+1} = x_n + y_n + \epsilon_{x,n} \mod 2\pi$$

 $y_{n+1} = (1 - \nu)y_n + f_0 \sin x_{n+1} + \epsilon_{y,n}$

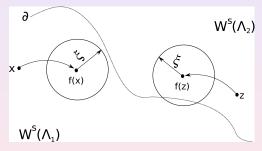
with iid random variables $\epsilon_n = (\epsilon_{x,n}, \epsilon_{y,n})$ drawn from uniform distribution bounded by $||\epsilon_n|| < \xi$ of noise amplitude ξ perturbed dynamics $\mathbf{F}(\mathbf{x}_i) = \mathbf{f}(\mathbf{x}_i) + \epsilon_i$:



From attractors to hopping on pseudo attractors

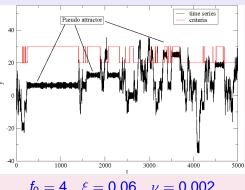
Consequences of the random perturbations:

• beyond a noise threshold $\xi \geq \xi_0$ the attracting sets $W^S(\Lambda_i)$ lose their stability due to holes



- the (invariant) attractors become (quasi-invariant) pseudo attractors from which there is noise-induced escape
- the noise induces a hopping process between all coexisting pseudo attractors

the resulting perturbed dissipative dynamics is intermittent:



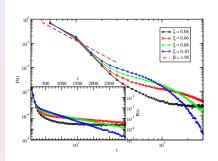
$$f_0 = 4$$
, $\xi = 0.06$, $\nu = 0.002$

 stickiness to pseudo attractors measured by criterion that maximal eigenvalue of the Jacobian matrix along orbit < 1

Escape time distribution

Outline

probability distributions P(t) of escape times t from pseudo attractors computed by using eigenvalue criterion (plus a Markov assumption and averaging over all non-uniform pseudo attractors):



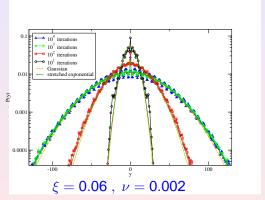
dissipation $\nu = 0.002$ with different noise strength ξ

- transition from power law (stickiness) to exponential
- transition takes longer when $\xi \to 0$

Diffusion

Outline

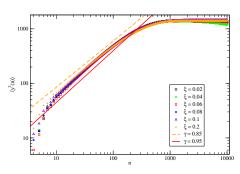
probability distribution function $P_n(y)$ for position y at different time steps n:



- there is Gaussian-like diffusive spreading up to n < 1000
- localization trivially due to boundedness of pseudo attractors

Summary

mean square displacement $< y^2(n) >$ for position y and different noise amplitudes ξ at $\nu = 0.002$:



- transient subdiffusion $< y^2(n) > \sim n^{\gamma}$ up to n < 1000
- only small variation of the subdiffusive exponent $0.85 < \gamma < 0.95$ for different ξ

Summary

Continuous time random walk theory

reproduce simulation results by CTRW theory (Montroll, Weiss, Scher, 1973): define stochastic process by master equation with waiting time distribution w(t) and jump distribution $\lambda(x)$

$$\varrho(\mathbf{x},t) = \int_{-\infty}^{\infty} d\mathbf{x}' \lambda(\mathbf{x} - \mathbf{x}') \int_{0}^{t} dt' \ w(t - t') \ \varrho(\mathbf{x}',t') +$$
$$+ (1 - \int_{0}^{t} dt' \ w(t')) \delta(\mathbf{x})$$

structure: jump + no jump for points starting at (x, t) = (0, 0)Fourier-Laplace transform yields Montroll-Weiss eqn (1965)

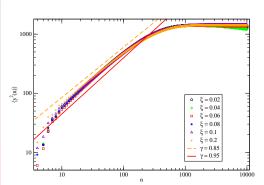
$$\hat{\varrho}(k,s) = \frac{1 - \tilde{w}(s)}{s} \frac{1}{1 - \hat{\lambda}(k)\tilde{w}(s)}$$

with mean square displacement $\langle x^2(s) \rangle = -\frac{\partial^2 \tilde{\varrho}(k,s)}{\partial k^2}$

Motivation

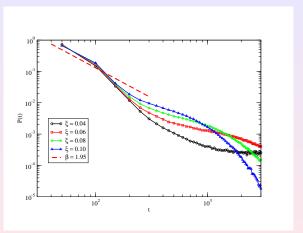
CTRW theory and mean square displacement

CTRW theory predicts that solving the MW eqn. for a power law waiting time distribution $w(t) \sim t^{-(\gamma+1)}$ with jump distribution $\lambda(x) = \delta(|x| - const.)$ yields $< x^2(t) > \sim t^{\gamma}$



for $\nu = 0.002$, $\xi = 0.06$ we have $\langle y_n^2 \rangle \sim n^{\gamma}$ with $\gamma \simeq 0.95$

CTRW theory and escape time distribution



the dashed red line represents the CTRW theory prediction of $P(t) \sim t^{-1.95}$ corresponding to $< v^2(n) > \sim n^{0.95}$

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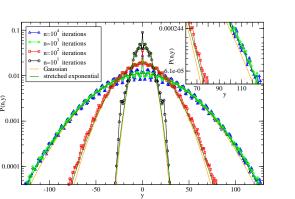
CTRW theory

CTRW theory and position pdf

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CTRW theory also predicts a stretched exponential position pdf, here: $P_n(y) \sim \exp(-cx^{2/(2-\gamma)})$



green lines represent the CTRW theory pdf for $\gamma = 0.95$: corrects the mismatch to Gaussian in the tails

CTRW theory

Summary

- central theme: study of diffusion generated by randomly perturbing dissipative deterministic dynamics
- main result: for the dissipative standard map non-hyperbolic stickiness to pseudo attractors under random perturbations generates
 - power law escape time distributions and
 - stretched exponential position distributions leading to
 - subdiffusion

simulation results consistently explained by CTRW theory

 outlook: similar phenomena in other randomly perturbed deterministic dynamical systems?

reference:

C.S.Rodrigues A.V.Chechkin, A.P.S. de Moura, C.Grebogi, R.Klages, Europhys.Lett. **108**, 40002 (2014)